

Lecture II: Nonlinear Schrödinger Equation

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Recall: our goal is to solve the (NLS)

$$\begin{cases} i\partial_t u + \Delta u = |u|^2 u \\ u|_{t=0} = u_0(x) \end{cases}$$

(NLS) is a Hamiltonian PDE with two conservation laws:

mass: $M[u] := \|u\|_{L^2}^2 = \int_{\mathbb{T}^2} |u|^2 dx$

energy: $H[u] := \int_{\mathbb{T}^2} \frac{1}{2} |\nabla u|^2 + \frac{1}{4} |u|^4 dx$

proposition: If u is a smooth solution on \mathbb{T}^2 , then

(Exercise) $\frac{d}{dt} M[u] = 0 = \frac{d}{dt} H[u(t)]$

II.1): Sobolev spaces on \mathbb{T}^d .

guiding principle: Define some norms for smooth functions and then consider the completion w.r.t. this norm.

$$\forall m \in \mathbb{N}. \quad \|u\|_{H^m(\mathbb{T}^d)}^2 = \sum_{|\alpha| \leq m} \|\partial^\alpha u\|_{L^2(\mathbb{T}^d)}^2$$

notation: $\partial^\alpha = \partial_{x_1}^{\alpha_1} \cdots \partial_{x_d}^{\alpha_d}$, $\alpha = (\alpha_1, \dots, \alpha_d) \in \mathbb{N}^d$.

$$|\alpha| = \alpha_1 + \dots + \alpha_d.$$

Definition: $H^m(\mathbb{T}^d) = \{u \in C^\infty(\mathbb{T}^d)\}^{H^m} \rightarrow$ completion

Observation: $\widehat{\partial^\alpha u}(k) = (2\pi i k)^\alpha \widehat{u}(k)$

where $(2\pi i k)^\alpha = k_1^{\alpha_1} \cdots k_d^{\alpha_d} (2\pi i)^{|\alpha|}$

Example: $\widehat{\nabla u}(k) = 2\pi i k \widehat{u}(k)$

$$\widehat{\Delta u}(k) = -4\pi^2 |k|^2 \widehat{u}(k)$$

By Parseval Identity: $\|\partial^\alpha u\|_{L^2(\mathbb{T}^d)}^2 = \sum_{k \in \mathbb{Z}^d} |(2\pi i k)^\alpha \widehat{u}(k)|^2$

$$\Rightarrow \|u\|_{H^m(\mathbb{T}^d)}^2 \simeq \sum_{|\alpha| \leq m} \|k^\alpha \widehat{u}(k)\|_{L^2(\mathbb{Z}^d)}^2 \simeq \sum_{k \in \mathbb{Z}^d} |(1+|k|^m) \widehat{u}(k)|^2$$

For $s > 0$ (not necessarily an integer)

$$\|u\|_{H^s(\mathbb{T}^d)}^2 := \sum_{k \in \mathbb{Z}^d} |(1+|k|^s) \hat{u}(k)|^2$$

Proposition : $H^{s_1}(\mathbb{T}^d) \subseteq H^{s_2}(\mathbb{T}^d)$, $s_1 \geq s_2$

In particular, $H^s(\mathbb{T}^d) \subseteq L^2(\mathbb{T}^d)$, $s \geq 0$

Remark : $\|u\|_{H^{s_2}} \leq \|u\|_{H^{s_1}}$

Theorem (Sobolev embedding)

(i) $0 < s < \frac{d}{2}$, $H^s(\mathbb{T}^d) \subseteq L^p(\mathbb{T}^d)$, $p = \frac{2d}{d-2s}$

(ii) $s > \frac{d}{2}$, $H^s(\mathbb{T}^d) \subseteq L^\infty(\mathbb{T}^d)$ (not correct when $s = \frac{d}{2}$)

(iii) product rule: $s > \frac{d}{2}$, $\|f \cdot g\|_{H^s} \lesssim \|f\|_{H^s} \|g\|_{H^s}$.

example: $d=2$, $H^2(\mathbb{T}^d) \subseteq L^\infty(\mathbb{T}^d)$

pf:

(ii) it suffices to prove that $\forall f \in C^\infty(\mathbb{T}^d)$

(*) $\|f\|_{L^\infty(\mathbb{T}^d)} \leq C \|f\|_{H^s(\mathbb{T}^d)}$

$\forall x$, $f(x) = \sum_{k \in \mathbb{Z}^d} \hat{f}(k) e^{2\pi i k x}$

$$\Rightarrow |f(x)| \leq \sum_{k \in \mathbb{Z}^d} |\hat{f}(k)| \underbrace{|e^{2\pi i k x}|}_{\leq 1}$$

$$\leq \sum_{k \in \mathbb{Z}^d} |\hat{f}(k)|$$

$$= \sum_{k \in \mathbb{Z}^d} |\hat{f}(k)| (1+|k|^s) \frac{1}{1+|k|^s}$$

$$\leq \left(\sum_{k \in \mathbb{Z}^d} |\hat{f}(k)|^2 (1+|k|^s)^2 \right)^{1/2} \left(\sum_{k \in \mathbb{Z}^d} \frac{1}{(1+|k|^s)^2} \right)^{1/2}$$

$$\left(\sum_{k \in \mathbb{Z}^d} \frac{1}{(1+|k|^s)^2} \right)^{1/2} \sim \int_{\mathbb{R}^d} \frac{1}{(1+|x|^s)^2} dx < \infty \quad , \quad s > \frac{d}{2}$$

So $\|f\|_{L^\infty(\mathbb{T}^d)} \leq C \|f\|_{H^s(\mathbb{T}^d)}$.

(iii) Consider the Fourier transform

$$(\widehat{f * g})(k) = \sum_{m \in \mathbb{Z}^d} \widehat{f}(m) \widehat{g}(k-m)$$

$$\widehat{f * g} = \widehat{f} \widehat{g}$$

notation : $\langle k \rangle = \sqrt{1 + |k|^2} \simeq 1 + |k|$

$$\Rightarrow \|f * g\|_{H^s}^2 = \sum_{k \in \mathbb{Z}^d} |k|^{2s} \left| \sum_m \widehat{f}(m) \widehat{g}(k-m) \right|^2$$

$$= \sum_{k \in \mathbb{Z}^d} |k|^{2s} \left| \sum_m \frac{\widehat{f}(m) \langle m \rangle^s}{\langle m \rangle^s} \frac{\widehat{g}(k-m) \langle k-m \rangle^s}{\langle k-m \rangle^s} \right|^2$$

note: $a_m = \widehat{f}(m) \langle m \rangle^s$, $b_m = \widehat{g}(m) \langle m \rangle^s \rightsquigarrow \|f\|_{H^s}^2 \simeq \sum a_m^2$

we observe that $\langle k \rangle \leq \langle k-m \rangle + \langle m \rangle$

$$|k| \leq |k-m| + |m|$$

so. $\langle k \rangle^s \leq C_s (\langle k-m \rangle^s + \langle m \rangle^s)$

$$\Rightarrow \|f * g\|_{H^s}^2 \leq C_s \sum_{k \in \mathbb{Z}^d} \left| \sum_{m \in \mathbb{Z}^d} (\langle k-m \rangle^s + \langle m \rangle^s) \frac{a_m b_{k-m}}{\langle m \rangle^s \langle k-m \rangle^s} \right|^2$$

$$\leq C_s \sum_k \left| \sum_m \frac{a_m b_{k-m}}{\langle m \rangle^s} + \frac{a_m b_{k-m}}{\langle k-m \rangle^s} \right|^2$$

$$\leq C_s \left[\sum_k \left| \sum_m \frac{a_m b_{k-m}}{\langle m \rangle^s} \right|^2 + \sum_k \left| \sum_m \frac{a_m b_{k-m}}{\langle k-m \rangle^s} \right|^2 \right]$$

$$\leq C_s \left[\sum_k \sum_m a_m^2 b_{k-m}^2 \underbrace{\sum \frac{1}{\langle m \rangle^{2s}}}_{\leq C \text{ if } s > \frac{d}{2}} + \sum_k \sum_m a_m^2 b_{k-m}^2 \underbrace{\sum \frac{1}{\langle k-m \rangle^{2s}}}_{\leq C \text{ if } s > \frac{d}{2}} \right]$$

$$\leq C_s \|f\|_{H^s}^2 \|g\|_{H^s}^2$$

□

II. 2): Local well-posedness of (NLS) with H^2 -data. [P4]

$$(NLS) : \begin{cases} i \partial_t u + \Delta u = |u|^2 u. & (t, x) \in \mathbb{R} \times \mathbb{T}^2 \\ u|_{t=0} = u_0 \in H^2(\mathbb{T}^2) \end{cases}$$

① Consider Linear Schrödinger Equation:

$$LS \quad \begin{cases} i \partial_t u + \Delta u = 0 \\ u|_{t=0} = u_0 \end{cases}$$

$$\Rightarrow i \partial_t \hat{u}(t, k) - 4\pi^2 |k|^2 \hat{u}(t, k) = 0$$

$$\Rightarrow \partial_t \hat{u}(t, k) + 4\pi^2 i |k|^2 \hat{u}(t, k) = 0$$

$$\Rightarrow \hat{u}_k(t, k) = e^{-4\pi^2 i |k|^2 t} \hat{u}_0(k)$$

notation: $u_{li}(t, x) = e^{it\Delta} u(0, x)$

$$f(t, x) = \sum_{k \in \mathbb{Z}^2} \hat{f}(t, k) e^{2\pi i k x}$$

$$\Rightarrow u_{li}(t, x) = \sum_{k \in \mathbb{Z}^2} e^{-4\pi^2 i |k|^2 t + 2\pi i k x} \hat{u}_0(k)$$

② Inhomogeneous case.

$$\begin{cases} i \partial_t v + \Delta v = F \\ v|_{t=0} = v_0 \end{cases}$$

$$\Rightarrow v(t, \cdot) = e^{it\Delta} v_0 + \frac{1}{i} \int_0^t e^{i(t-s)\Delta} F(s, \cdot) ds$$

Now, rewrite (NLS) as an integral equation:

$$(*) \quad u(t) = e^{it\Delta} u_0 + \frac{1}{i} \int_0^t e^{i(t-s)\Delta} (|u|^2 u(s)) ds \quad (\text{Duhamel formula})$$

(check it by expanding as Fourier series)

Theorem (Local well-posedness)

Given $u_0 \in H^2(\mathbb{T}^2)$, there exists a unique solution $u(t) \in C([0, T_{\max}), H^4)$

of (NLS) in the sense:

$$u(t) = e^{it\Delta} u_0 + \frac{1}{i} \int_0^t e^{i(t-s)\Delta} (|u|^2 u(s)) ds$$

The maximal lifespan $T_{\max} \gtrsim \frac{1}{\|u_0\|_{H^2}^3}$

Moreover, we have the blow-up criteria:

(P5)

$$(*) \quad \limsup_{t \nearrow T_{\max}} \|u(t)\|_{H^2} = +\infty$$

Tool (1) Fixed point argument: consider the Cauchy problem

$$(Eq) \quad \partial_t u(t) = F(u(t)) \text{ for all } t \in I, \quad u(t_0) = u_0$$

where the interval I , the initial time t_0 , the initial datum $u_0 \in \mathcal{D}$, and the nonlinearity $F: \mathcal{D} \rightarrow \mathcal{D}$ are given. ↑
normed
vector space

A strong solution of (Eq) is a function $u \in C_{loc}^0(I \rightarrow \mathcal{D})$ which solves

$$(Eq) \text{ in the integral sense that } u(t) = u_0 + \int_{t_0}^t F(u(s)) ds.$$

Theorem: Let $F \in C^{0,1}(\mathcal{D} \rightarrow \mathcal{D})$ be a Lipschitz function on \mathcal{D} with Lipschitz constant $\|F\|_{C^{0,1}} = \mu$. Let $0 < T < 1/\mu$. Then for any $t_0 \in \mathbb{R}$ and $u_0 \in \mathcal{D}$, there exists a (strong) solution $u: I \rightarrow \mathcal{D}$ to the Cauchy pb (Eq) where $I = [t_0 - T, t_0 + T]$.

(2) the way to prove well-posedness, at least locally, is by defining an operator

$$Lu = S(t)u_0 + c \int_0^t S(t-s) (|u|^2 u(s)) ds$$

and then showing that in a certain space of functions X one has a fixed point and as a consequence a solution according to (*).

The hard part is to decide what space X could work.

Global well-posedness

Thanks to blow-up criteria, to show that solutions are indeed global. (i.e. $T_{\max} = +\infty$), we have to prove that $\|u(t)\|_{H^2}$ is finite for any

finite t , provided that the solution exists.

Remark: In principal, to control $\|u(t)\|_{H^s}$ global in time, P6

we need to search for some Lyapunov functional, that controls the H^s -norm. For example, the energy conservation law controls $\|u(t)\|_{H^1}$ globally in time. However, we don't have such conserved quantity at level H^2 .

Fortunately, for the specific cubic NLS equation in 2D, we are able to adapt a Gronwall type argument to bound $\|u(t)\|_{H^2}$ with a bad control in time.

A key tool is the following Brezis-Gallouet's inequality.

Lemma (Brezis-Gallouet)

Let $s > \frac{d}{2}$. Then for any $u \in H^s(\mathbb{T}^d)$.

$$\|u\|_{L^\infty} \lesssim \|u\|_{H^1} \left[1 + \log^{1/2} \left(1 + \frac{\|u\|_{H^s}}{\|u\|_{H^1}} \right) \right]$$

proposition: $\forall u_0 \in H^2(\mathbb{T}^2)$. the solution $u(t)$ of (NLS) is global.

In particular, $\|u(t)\|_{H^2} \lesssim C(\|u_0\|_{H^2}) e^{ct}$.

pf: Recall (NLS) $i\partial_t u + \Delta u = |u|^2 u$

what is $\|u(t)\|_{H^2}$? $\|u\|_{H^2}^2 \simeq \|\Delta u\|_{L^2}^2 + \|u\|_{L^2}^2$

\Rightarrow Do some estimate for Δu

$$\frac{d}{dt} \|\Delta u\|_{L^2(\mathbb{T}^2)}^2 = \frac{d}{dt} \int_{\mathbb{T}^2} \Delta u \Delta \bar{u} \, dx$$

$$= 2 \operatorname{Re} \int_{\mathbb{T}^2} \Delta (u_t) \Delta \bar{u} \, dx$$

$$i u_t = -\Delta u + |u|^2 u$$

$$= 2 \operatorname{Re} \int_{\mathbb{T}^2} \Delta (i \Delta u - i |u|^2 u) \Delta \bar{u} \, dx$$

$$u_t = i \Delta u - i |u|^2 u$$

$$= 2 \operatorname{Re} \left(i \int_{\mathbb{T}^2} \Delta (\Delta u) \Delta \bar{u} \, dx \right) - 2 \operatorname{Re} \left(i \int_{\mathbb{T}^2} \Delta (|u|^2 u) \Delta \bar{u} \, dx \right)$$

IPP

$$= -2 \operatorname{Re} \left(i \int_{\mathbb{T}^2} |\nabla \Delta u|^2 \, dx \right) - 2 \operatorname{Re} \left(i \int_{\mathbb{T}^2} \Delta (|u|^2 u) \Delta \bar{u} \, dx \right)$$

$$\leq 2 \|\Delta(u^2 u)\|_{L^2(\mathbb{T}^2)} \|\Delta u\|_{L^2(\mathbb{T}^2)}$$

$$\leq 2 \|u\|_{H^2(\mathbb{T}^2)} \|u\|_{L^\infty(\mathbb{T}^2)}^2 \|\Delta u\|_{L^2(\mathbb{T}^2)}$$

$$\Rightarrow \frac{d}{dt} \|u\|_{H^2(\mathbb{T}^2)}^2 \leq C \|u\|_{H^2(\mathbb{T}^2)}^2 \|u\|_{L^\infty(\mathbb{T}^2)}^2$$

By Brezis-Gallouet,

$$\|u\|_{L^\infty}^2 \lesssim \|u\|_{H^1}^2 \left[1 + \log^{1/2} \left(1 + \frac{\|u\|_{H^2}}{\|u\|_{H^1}} \right) \right]^2$$

we have

$$\frac{d}{dt} \|u\|_{H^2(\mathbb{T}^2)}^2 \leq C \|u\|_{H^2(\mathbb{T}^2)}^2 \|u\|_{H^1(\mathbb{T}^2)}^2 \left[1 + \log \left(1 + \frac{\|u\|_{H^2}}{\|u\|_{H^1}} \right) \right]$$

By conservation law. $\frac{d}{dt} H[u(t)] = 0$, $H[u(t)] = \int_{\mathbb{T}^2} \frac{1}{2} |\nabla u|^2 + \frac{1}{4} |u|^4 dx$

$$\Rightarrow \|u(t)\|_{H^1(\mathbb{T}^2)} \leq C(\|u_0\|_{H^1(\mathbb{T}^2)})$$

So

$$\frac{d}{dt} \|u\|_{H^2(\mathbb{T}^2)}^2 \leq C(\|u_0\|_{H^1(\mathbb{T}^2)}) \|u\|_{H^2(\mathbb{T}^2)}^2 \left[1 + \log(1 + \|u\|_{H^2}) \right]$$

we arrive at an inequality of form

$$f'(t) \leq C f(t) [1 + \log(1 + f(t))]$$

$$\Rightarrow f(t) \leq C e^{ct}$$

Gronwall

□.

Pf of Brezis-Gallouet lemma:

$$\|u\|_{L^\infty} \leq \|\hat{u}(k)\|_{l_k^1}$$

$$= \underbrace{\sum_{|k| \leq \lambda} |\hat{u}(k)|}_{\text{I}} + \underbrace{\sum_{|k| > \lambda} |\hat{u}(k)|}_{\text{II}}$$

$$\text{I} = \sum_{|k| \leq \lambda} |\hat{u}(k)| (1 + |k|) \frac{1}{1 + |k|}$$

$$\leq \|u\|_{H^1(\mathbb{T}^2)} \left(\sum_{|k| \leq \lambda} \frac{1}{(1 + |k|)^2} \right)^{1/2}$$

$$\leq C (\log \lambda)^{1/2} \|u\|_{H^1(\mathbb{T}^2)}$$

$$\text{II} = \sum_{|k| > \lambda} |\hat{u}(k)| (1 + |k|^2) \frac{1}{1 + |k|^2}$$

$$\leq \|u\|_{H^2} \left(\sum_{|k| > \lambda} \frac{1}{(1+|k|^2)^2} \right)^{1/2}$$

$$\leq \|u\|_{H^2} \left(\sum_{|k| > \lambda} \frac{1}{|k|^4} \right)^{1/2}$$

$$\leq C \frac{1}{\lambda} \|u\|_{H^2(\mathbb{T}^2)}$$

$$\Rightarrow \|u\|_{L^\infty} \leq C \left(\|u\|_{H^1} (\log \lambda)^{1/2} + \frac{\|u\|_{H^2}}{\lambda} \right)$$

$$\text{Let } \lambda = 1 + \frac{\|u\|_{H^2}}{\|u\|_{H^1}} \Rightarrow \frac{\|u\|_{H^2}}{\lambda} \leq \|u\|_{H^1}$$

$$\Rightarrow \|u\|_{L^\infty} \leq C \|u\|_{H^1} \log^{1/2} \left(1 + \frac{\|u\|_{H^2}}{\|u\|_{H^1}} \right)$$

□.